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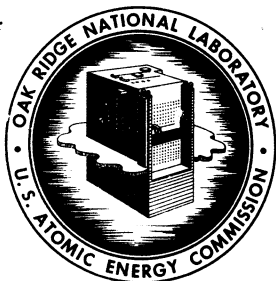
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DATE: November 12, 1962
SUBJECT: Preliminary Equations to Describe Iodine and Xenon
Behavior in the MSRE
TO: Distribution
FROM: J. R. Engel

ABSTRACT

Equations are presented to describe the behavior of iodine and xenon in the primary loop of the MSRE at steady state. These equations may be programmed for machine solution to obtain the spatial distribution of xenon in the reactor from which the xenon poison fraction can be estimated. An alternate use of the equations is to evaluate physical parameters from observations made on the operating reactor.

An attempt has been made to include all of the important, or potentially important, behavior mechanisms in the equations. Comments and suggestions are solicited regarding the mechanisms treated, other possible mechanisms, and the mathematical representations thereof.

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INTRODUCTION

Xenon poisoning is a problem of significant interest in the MSRE. A reasonably accurate evaluation of the xenon poison fraction requires a detailed description of the xenon behavior in the entire core loop. Since most of the Xe is formed by decay of iodine, an equally detailed description of the iodine behavior is also required. This memo presents a set of steady-state equations which is intended to fulfill these requirements. A number of effects which may or may not be important in the reactor are included while some, for which there is good evidence that they are unimportant, are omitted. The equations describe the fission product behavior in terms of a number of physical properties of the system, the fuel and the fission products, even though it may not be possible at this time to assign accurate values to all of the properties.

This set of equations may be used in two ways. First, if values are assigned to all of the physical parameters, the equations can provide an estimate of the xenon distribution and, hence, the poison fraction in the reactor. Second, the equations may be used with the operating reactor to evaluate those physical parameters which cannot otherwise be evaluated. For the latter case it is hoped that the reactor can be operated under a sufficiently wide variety of conditions to permit simultaneous evaluation of a number of the parameters. The use of time-dependent forms of the equations and observation of xenon transients may be used to aid these evaluations.

The purpose of this memo is to solicit comments and suggestions regarding the mechanisms and their mathematical representations. To this end the text describes all the assumptions and approximations that were used in developing the equations as well as the equations themselves. It is expected that revisions to the equations will be required after all comments are received. The final equations will then be programmed for solution on the IBM-7090 computer to predict the xenon poison fraction. Ultimately, a program will be prepared to evaluate physical parameters from information obtained from the operating reactor.

It is anticipated that the equations which are initially programmed for computer solution will be simplified approximations of the ones presented here. Part of this simplification can be achieved by ignoring

behavior mechanisms of questionable importance. However, as a first step it seems desirable to include all mechanisms which might be important. Then, if the simple approximation proves to be inadequate for evaluating the reactor performance, a more refined analysis can be made omitting the simplifications.

SUMMARY OF BEHAVIOR MECHANISMS

Several mechanisms are considered for the appearance and disappearance of Xe and I in various parts of the reactor system. These mechanisms may be summarized as follows:

Appearance of Xe in circulating fuel

1. Direct production by fission of uranium in the circulating fuel
2. Decay of I in the circulating fuel
3. Desorption from metal walls

Disappearance of Xe from circulating fuel

1. Radioactive decay
2. Neutron absorption
3. Stripping in the fuel-pump bowl
4. Sorption in the graphite

Appearance of Xe in graphite

1. Direct production by fission of uranium in fuel soaked into the graphite
2. Decay of I in the graphite
3. Sorption from the circulating fuel

Disappearance of Xe from graphite

1. Radioactive decay
2. Neutron absorption

Appearance of Xe on metal walls

1. Decay of I sorbed on the walls

Disappearance of Xe from metal walls

1. Radioactive decay
2. Desorption into the circulating fuel

Appearance of I in circulating fuel

1. Direct production by fission of uranium in the circulating fuel

Disappearance of I from circulating fuel

1. Radioactive decay
2. Sorption in the graphite
3. Sorption on metal walls

Appearance of I in graphite

1. Direct production by fission of uranium in fuel soaked into the graphite
2. Sorption from the circulating fuel

Disappearance of I in graphite

1. Radioactive decay

Appearance of I on metal walls

1. Sorption from the circulating fuel

Disappearance of I from metal walls

1. Radioactive decay

It may be noted that all mechanisms are not applied equally to both species. In the case of the sorption processes, only the net currents are represented and the reverse processes of those given are taken into account in the concentration gradients. Other mechanisms, neutron absorption in I and I-stripping, are omitted because they are expected to be negligibly small.

DESCRIPTION OF EQUATIONS

The equations describing the behavior of iodine and xenon in the reactor were derived from material balance considerations. In developing these equations, the reactor system was divided into two parts: a core in which the detailed spatial distribution of xenon in both the fuel and graphite was of primary interest and an external loop where the spatial distribution was less important. With this goal in mind, the steady-state equations for the core were written in differential

form while simplified, integrated forms of the equations were used for the external loop.

Core Equations

The MSRE core consists of fuel salt flowing through channels and an array of graphite stringers; both of these materials were treated separately. The flow pattern of the fuel makes a further division in the mathematical treatment necessary. Although the established fluid flow in most of the core channels is laminar, there are regions of high turbulence at the channel inlets. In addition, turbulent flow may persist over the entire length of some of the channels. Two sets of fuel equations were developed to cover the different flow patterns. In solving the xenon distribution problem, the fuel channels will be treated in two parts, using the turbulent-flow equations for the lower parts and the laminar-flow equations for the remainder. Although the transition to laminar flow does not occur at a sharp boundary, it is expected that a reasonable approximation to the true condition can be obtained by judicious selection of a fixed transition elevation.

Fuel in Laminar Flow

Under laminar flow conditions it is necessary to consider variations in concentration in both the radial and axial directions within individual channels. To simplify the treatment the channels are regarded as circular with a diameter equal to the hydraulic diameter of the MSRE channels and the neutron flux is assumed constant in the transverse direction in any given channel. Material balances may then be written for the volume element,¹ $2\pi r' dr' dz$. For xenon

$$\begin{aligned}
 & Y_{xe} \sum_f^l \phi(r, z) r' dr' dz - \lambda_{xe} N_{xe}^l(r', z) r' dr' dz + \lambda_{I-I} N_{I-I}^l(r', z) r' dr' dz \\
 & + D_{xe}^l \nabla^2 N_{xe}^l(r', z) r' dr' dz - \sigma_{xe}^a \phi(r, z) N_{xe}^l(r', z) r' dr' dz \\
 & - v(r, r') \frac{\partial N_{xe}^l(r', z)}{\partial z} r' dr' dz = 0 \quad . \quad (1)
 \end{aligned}$$

¹See Appendix for definitions of symbols.

The various terms represent differential changes in the xenon population in the volume element due to the following mechanisms, in the order of the terms: (1) direct production from fission; (2) loss by radioactive decay; (3) production by decay of iodine; (4) net diffusion from the volume element; (5) burnup; and (6) net transport by the flowing fluid.

The diffusion term, as written, includes diffusion in both the radial and axial directions. In the actual reactor it is likely that axial diffusion will be much less important than axial transport by the flowing stream. If this is assumed the diffusion term may be simplified as follows:

$$D_{\text{Xe}}^{\ell} \nabla^2 N_{\text{Xe}}^{\ell}(r',z) r' dr' dz = D_{\text{Xe}}^{\ell} \left[\frac{\partial^2 N_{\text{Xe}}^{\ell}(r',z)}{\partial r'^2} + \frac{1}{r'} \frac{\partial N_{\text{Xe}}^{\ell}(r',z)}{\partial r'} \right] r' dr' dz \quad (2)$$

The microscopic velocity, $v(r')$, varies about a mean, or macroscopic, velocity for the channel. In addition, the macroscopic velocity varies with the radial position of the channel in the reactor. Axial variations in velocity, both macroscopic and microscopic, are neglected.

An equation similar to (1) can be written for iodine except that fewer terms are required because fewer mechanisms are involved in the iodine behavior. Thus:

$$Y_{\text{I}} \int_f^{\ell} \phi(r,z) r' dr' dz - \lambda_{\text{I}} N_{\text{I}}^{\ell}(r',z) r' dr' dz + D_{\text{I}}^{\ell} \nabla^2 N_{\text{I}}^{\ell}(r',z) r' dr' dz - v(r,r') \frac{\partial N_{\text{I}}^{\ell}(r',z)}{\partial z} r' dr' dz = 0 \quad (3)$$

All of the comments on the xenon equation, including an expression similar to (2) apply to the iodine equation.

Fuel in Turbulent Flow

The microscopic radial variations in concentration and velocity disappear in those channels, or parts of channels, where the fluid flow is turbulent. As a result the diffusion terms do not appear in the material balance equations for this condition. However, a new term must be included to describe the transport of xenon (or iodine) through the fluid film, at the channel wall, to the graphite. The volume element for this condition is $\pi R'^2 dz$. Then, for xenon

$$Y_{xe} \sum_f^l \phi(r,z) dz - \lambda_{xe} N_{xe}^l(z) dz + \lambda_I N_I^l(z) dz - \frac{2}{R} h_{xe} * \\ \left[N_{xe}^l(z) - N_{xe}^l(R',z) \right] dz - \sigma_{xe}^a \phi(r,z) - v(r) \frac{\partial N_{xe}^l(z)}{\partial z} dz = 0 \quad (4)$$

The quantity, $N_{xe}^l(R',z)$, is the concentration of xenon in the liquid on the graphite side of the fluid film and $N_{xe}^l(z)$ is the bulk concentration in the liquid. The difference between these quantities is the driving force for mass transfer across the film.

The equation for iodine which corresponds to (4) is

$$Y_I \sum_f^l \phi(r,z) dz - \lambda_I N_I^l(z) - \frac{2}{R} h_I \left[N_I^l(z) - N_I^l(R',z) \right] dz \\ - v(r) \frac{\partial N_I^l(z)}{\partial z} dz = 0 \quad (5)$$

Graphite

In order to simplify the mathematics somewhat, the graphite stringers are regarded as cylinders with an equivalent radius, R'' , to be assigned on the basis of the physical dimensions. This treatment is equivalent to regarding the reactor as an array of graphite and fuel cylinders in a medium which has infinite resistance to mass transfer in the axial direction and zero resistance in the radial direction. The material balance equations for the graphite are written on the assumption that, if fuel soaks into the graphite (providing a direct source of fission products), it will be uniformly distributed in all respects. As with the fuel channels, a uniform transverse neutron flux is assumed for the individual stringers.

For a graphite volume element, $2\pi r'' dr'' dz$, the material balance for xenon gives

$$Y_{xe} \sum_f^g \phi(r,z) r'' dr'' dz - \lambda_{xe} N_{xe}^g(r'',z) r'' dr'' dz + \lambda_I N_I^g(r'',z) r'' dr'' dz \\ + D_{xe}^{gr} \left[\frac{\partial^2 N_{xe}^g(r'',z)}{\partial r''^2} + \frac{1}{r''} \frac{\partial N_{xe}^g(r'',z)}{\partial r''} \right] r'' dr'' dz + D_{xe}^{gz} \frac{\partial^2 N_{xe}^g(r'',z)}{\partial z^2} r'' dr'' dz \\ - \sigma_{xe}^a \phi(r,z) N_{xe}^g(r'',z) r'' dr'' dz = 0 \quad (6)$$

This equation is similar to (1) except that the transport term does not appear because the graphite is stationary. The diffusion in graphite is represented as two separate terms and different symbols are used for the diffusion coefficient in the terms; it has been suggested that the values of the diffusion coefficient may be significantly different for the two directions.²

The material balance for iodine in the graphite volume element gives

$$\begin{aligned}
 Y_I \sum_f^g \phi(r, z) r'' dr'' dz - \lambda_I N_I^g(r'', z) r'' dr'' dz \\
 + D_I^{gr} \left[\frac{\partial^2 N_I^g(r'', z)}{\partial r''^2} + \frac{1}{r''} \frac{\partial N_I^g(r'', z)}{\partial r''} \right] r'' dr'' dz \\
 + D_I^{gz} \frac{\partial^2 N_I^g(r'', z)}{\partial z^2} r'' dr'' dz \quad . \quad (7)
 \end{aligned}$$

As before, the comments on the xenon equation also apply for iodine.

Boundary Conditions

A number of boundary conditions are required to fill out the set of equations for the core. Some of these are intuitively obvious while others require assumptions about the behavior mechanisms. Some boundary conditions also depend on the nature of the flow in the fuel channels.

Fuel-Graphite Interface -- Two boundary conditions can be specified at the fuel-graphite interface. The first of these requires that the net currents in the two media at the interface be equal; the form of this expression depends on the type of fluid flow. For xenon, where the channel flow is laminar

$$D_{xe}^{gr} \frac{\partial N_{xe}^g(R'', z)}{\partial r''} = D_{xe}^l \frac{\partial N_{xe}^l(R', z)}{\partial r'} \quad . \quad (8)$$

For iodine at the same condition

$$D_I^{gr} \frac{\partial N_I^g(R'', z)}{\partial r''} = D_I^l \frac{\partial N_I^l(R', z)}{\partial r'} \quad . \quad (9)$$

²A. Taboada, personal communication